

MOLECULAR STRUCTURE, HOMO-LUMO AND VIBRATIONAL ANALYSIS OF ERGOLINE BY DENSITY FUNCTIONAL THEORY

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Abstract: In this work, quantum chemical study on a natural product ergoline has been presented using density functional theory (DFT) employing 6-311++G (d,p) basis set. A complete vibrational assignment has been performed for the theoretical FT-IR and Raman wavenumbers along with the potential energy distribution (PED) with the result of quantum chemical calculations. The structure–activity relationship has been interpreted by mapping electrostatic potential surface (MEP). Graphical representation of frontier molecular orbitals with their energy gap have been analyzed theoretically for both the gaseous and solvent environment employing time dependent density functional theory (TD-DFT) employing 6-31G basis set.

Keywords: Ergoline; *ab initio*; Vibrational spectroscopy; Density functional theory.

INTRODUCTION

The plant derived natural products cover a large field of medicine, which play very important role in the life threatening conditions due to their biological activity. Ergoline alkaloids (also called ergot alkaloids), the natural products occurring in the flowering plants *Rivea corymbosa* and *Ipomoea violacea* (Convolvulaceae; called morning glory) and taxonomically unrelated fungal, are known to inhibit the release of prolactin^{1,2}. Ergot alkaloids, the derivatives of the tetracyclic ergoline scaffold, divided into three main groups; display a diverse spectrum of pharmacological properties; including central, peripheral, and neurohormonal activities that arise from their activities at various monoaminergic receptors³. Naturally occurring alkaloids like; the D-lysergic acid amides, produced by the “ergot fungus” *Claviceps purpurea*, have been used as medicinal agents for a long time. They are active ingredients in medications designed to treat migraine disease or are used in childbirth and weaning⁴⁻⁶.

Ergoline alkaloids are also used as hallucinogens by native Central American people in religious ceremonial practices⁷. Historically, some ergots (e.g., ergoline and ergopeptine) have been used as probes to study various 5-HT receptor subtypes⁸⁻¹². However, during the middle Ages, they were known for poisoning and severe toxicosis in cattle feeding with ergoline alkaloid producing clavicipitaceous fungi¹³⁻¹⁶.

Alkaloids occupy unique place in the natural product. Their complex structures and dynamical behaviors attract the chemists towards the challenging problems in structural elucidation and synthesis. Vibrational spectroscopy is a valuable method for determining dynamical behavior and the electronic structure of the alkaloids microscopically¹⁷. Raman and infrared (IR) are the traditional methods of vibrational analysis for the structural characterization of the substances¹⁸. They have been widely used to study energetic, thermal, electronic and dynamical behavior from vibrational dynamics of the

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large number of biomolecules. In the recent years, the application of *ab initio* methods has also become a part of keen interest as they provide additional vibrational spectroscopic data^{17,19,20}. In this study, the theoretical vibrational spectra of ergoline (6aR)-4,6,6a,7,8,9,10a-Octahydroindolo[4,3-*fg*] quinoline) with the identification of various normal modes have been investigated. The aim of this study was to study the molecular structure, vibrational wavenumbers and the electronic absorption band theoretically. The molecular structure, harmonic vibrational wavenumbers, electrostatic potential surfaces, absolute Raman scattering activities and IR absorption intensities have been calculated by density functional theory (DFT)²¹ using Gaussian 09 program package²² employing B3LYP/6-311++G(d,p) basis set^{23,24}. The highest occupied molecular orbital (HOMO), lowest unoccupied molecular orbital (LUMO), energy band gap (ΔE) and the dipole moment were determined for the optimized structure. A theoretical time dependent density functional theory (TD-DFT)^{25,26} method has been used to calculate absorption parameters in gas phase employing 6-31G basis set. Graphical representation of HOMO and LUMO can provide valuable information that elucidates the nature of reactivity and some of the structural and physical properties of the molecule.

METHODOLOGY

Computational methods

The electronic structure and the optimized geometry of the molecule were computed by *ab initio* Hartree-Fock (HF) and the density functional functional theory (DFT) method using Gaussian 09 program²² employing 6-311++G(d,p) basis set. The DFT calculations were carried out in the frame-work of Becke-Lee-Yang-Parr [B3LYP] functional, in which the exchange functional is local spin density exchange with Becke gradient correlation²⁷ and the correlation functional is that of Lee, Yang and Parr with both local and non local terms^{28,29}. The molecular structure, vibrational frequencies and energy of optimized geometric structure of the molecule were calculated at

B3LYP/6-311++G (d,p) level using Gaussian 09 program²².

The absolute Raman scattering and IR absorption intensities are calculated in the harmonic approximation in the same level. The optimized ground state structure of ergoline is shown in figure 1. The normal mode analysis was performed and the potential energy distribution (PED) was calculated using localized symmetry using GAR2PED³⁰. For this purpose a complete set of 90 internal co-ordinates was defined by Pullay's recommendation^{31,32}. The graphical representation of MEP mapping and, the calculated IR and Raman spectra were made using Gauss View program³³. Visualization and confirmation of calculated data were done by using the CHEMCRAFT program³⁴. The frontier energy orbitals HOMO (characterized by highest occupied molecular orbital) and LUMO (characterized by lowest unoccupied molecular orbital), and band gap (ΔE) were calculated using TD-DFT^{25,26} employing 6-31G basis for both the gaseous and solvent phase.

RESULTS AND DISCUSSION

Geometry optimization

Using the standard geometric parameters³⁵, geometry optimization was performed as the first task in *ab initio* calculation without using any constraints. The optimized geometric parameters were used in the vibrational frequency calculations to characterize all the stationary points as minima. The related energies of the molecule are calculated employing *ab initio* HF functions and DFT functional. The energy calculated by DFT (- 652.8040 Hartree) is lower showing more stable than that by HF (- 648.5243 Hartree) due to some electron correlations in the former. The enthalpy difference between these two theories is 12.0725 kcal/mol. The optimized ground state structure (in the gas phase) is as shown in figure 1.

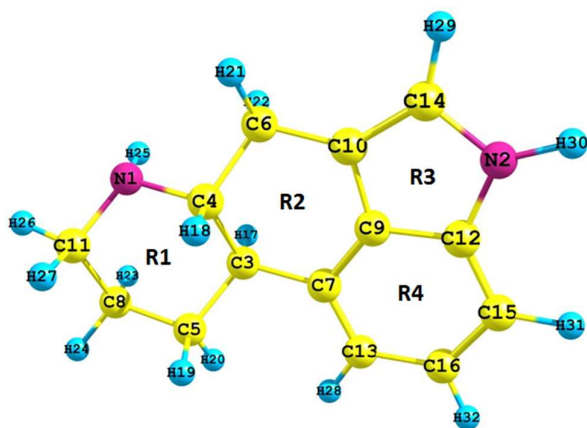


Figure 1: Optimized structure of ergoline.

Molecular electrostatic potential

The molecular electrostatic potential (MEP) in a molecule at a point $r(x,y,z)$ is the force on unitary positive test charge at that point due to its whole electrical charge and is given by:

$$V(r) = \sum_{A=1}^N \frac{Z_A}{|R_A - r|} - \int \frac{\rho(r') dr'}{|r' - r|}$$

where Z_A is the charge on nucleus A located at R_A and $\rho(r')$ is the electron density. The first term is due to the nucleus and the second due to electron cloud.

The MEP provides a visual method to understand the relative polarization of molecule¹⁷. Such surfaces depict the size, shape, charge density, and site of chemical reactivity of the molecules. In the surface generated, negative electrostatic potential (shades of red color) corresponds to an attraction of the proton by the concentrated electron density in the molecules (from lone pairs, pi-bonds, etc.) and positive electrostatic potential (shades of blue color) corresponds to repulsion of the proton by the atomic nuclei in the regions where low electron density exists and the nuclear charge is incompletely shielded. The largely white or lighter color shades on the surface indicate that the molecule is mostly non-polar. The potential increases in the order red < orange < yellow < green < blue.

The molecular electrostatic potential (MEP), electron

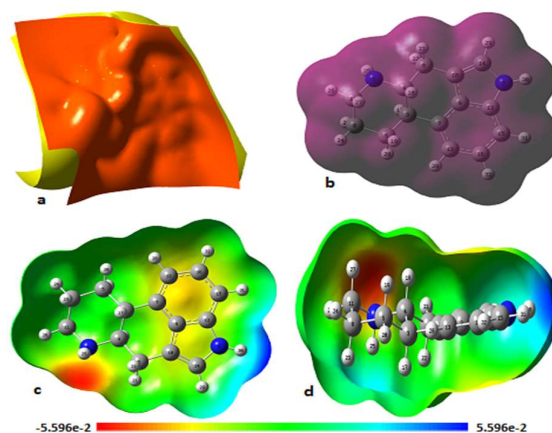


Figure 2: Graphical plots of: (a) electrostatic potential surface (b) electron density (c) molecular electrostatic potential mapped in-plane and (d) molecular electrostatic potential mapped perpendicular to the plane (from $-5.596e-2$ esu to $+5.596e-2$ esu).

density (ED) and the electrostatic potential surface (ESP) mapped with the output obtained by B3LYP/6-311++G (d,p) basis are shown in figures 2(a-d). From the figure, the highest negative potential (red region) is localized over the nitrogen atom N1. A less negative area occurs over the ring R2 and R3. The lowest potential with the blue region is localized over the hydrogen atom connected with N2 atom. The ED plot shows uniform charge distribution over the molecular surface.

UV-visible spectroscopy

UV-vis spectroscopy provides some valuable information about the nature of electronic transitions between the two prominent frontier molecular orbitals named as; HOMO (characterized by highest occupied molecular orbital) and LUMO (characterized by lowest unoccupied molecular orbital). The nature of electronic transition has been studied by using TD-DFT employing 6-31G basis set^{25,26} for both the gas and solvent (MeOH) phases. The theoretical absorption spectra are shown in the figure 3. The calculated electronic transitions with high oscillator strength are listed in Table 1.

Table 1. Electronic transitions with absorption wavelength (λ_{\max} in nm), excitation energy (eV), oscillator strength, frontier orbital energies (eV) and the dipole moment (μ in D).

| Excited states | Gas Phase | | | | MeOH Phase | | | | Transition type |
|----------------|------------------------|-------------------|------------------|---------------------|-----------------------|--------|------------------|---------------------|-------------------------|
| | Transitions | eV | λ_{\max} | Oscillator strength | Transitions | eV | λ_{\max} | Oscillator strength | |
| 1 | H \rightarrow L | 4.7240 | 262.45 | 0.1017 | H \rightarrow L | 4.6077 | 269.08 | 0.1338 | $\pi \rightarrow \pi^*$ |
| 2 | H -3 \rightarrow L | 6.1719 | 200.89 | 0.2617 | H \rightarrow L +1 | 5.8791 | 210.89 | 0.3184 | „ |
| 3 | H \rightarrow L+2 | 6.3689 | 194.67 | 0.1881 | H-1 \rightarrow L+1 | 6.0278 | 205.69 | 0.1682 | „ |
| 4 | H-1 \rightarrow L+2 | 6.5217 | 190.11 | 0.0798 | H-3 \rightarrow L | 0.1682 | 198.09 | 0.1968 | „ |
| 5 | H-2 \rightarrow L+1 | 6.7531 | 183.60 | 0.2367 | H-2 \rightarrow L+1 | 6.5887 | 188.18 | 0.3575 | .. |
| 6 | H-2 \rightarrow L+2 | 7.0700 | 175.37 | 0.1412 | H-2 \rightarrow L+2 | 6.9796 | 177.64 | 0.1235 | „ |
| 7 | H -3 \rightarrow L+1 | 7.5295 | 164.66 | 0.1117 | H-3 \rightarrow L+1 | 7.4289 | 166.89 | 0.1739 | „ |
| 8 | H -3 \rightarrow L+2 | 8.4539 | 146.66 | 0.1868 | H-3 \rightarrow L+2 | 8.3861 | 147.85 | 0.1552 | „ |
| | E_{HOMO} | E_{LUMO} | ΔE | μ (D) | | | | | |
| Gas | - 5.147798 | 0.035511 | 5.183310 | 1.3941 | | | | | |
| MetOH | - 5.277569 | - 0.162751 | 5.114818 | 1.8221 | | | | | |

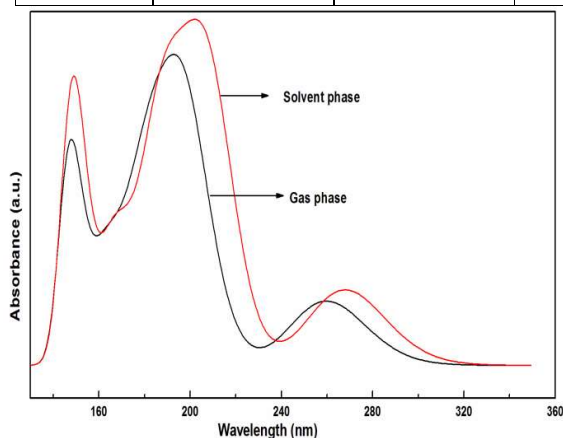
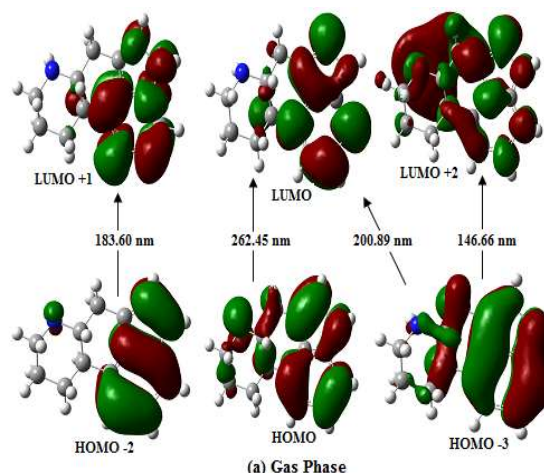


Figure 3: Plot of theoretical UV-vis spectra taken in gas (black) and the solvent (red).

HOMO and LUMO, the frontier molecular orbitals (FMOs), are very important quantum chemical parameters. These two vicinal orbitals play the same role of electron donor and acceptor, respectively. Their energy gap is an important stability index. A higher value of the energy difference more is the stability of the molecular system³⁶. The HOMO-LUMO plots are given in figure 4. TD-DFT calculations predict the intense electronic transitions at 262.45 nm (H \rightarrow L), 200.89 nm (H-3 \rightarrow L), 183.60 nm (H-

2 \rightarrow L+1) and 146.66 nm (H-3 \rightarrow L+2) with oscillator strengths 0.1017, 0.2617, 0.2367 and 0.1868, respectively in the gaseous phase. Similarly, in the solvent phase, the transitions are centered at 269.08 nm (H \rightarrow L), 205.69 nm (H-1 \rightarrow L+1), 210.89 nm (H \rightarrow L+1), 188.18 nm (H-2 \rightarrow L+1) and 147.85 nm (H-3 \rightarrow L+2) with the oscillator strengths 0.1338, 0.1682, 0.3184, 0.3575 and 0.1552, respectively. On the basis of calculated molecular orbital coefficient analysis, the main electronic excitations are due to $\pi \rightarrow \pi^*$ transitions.



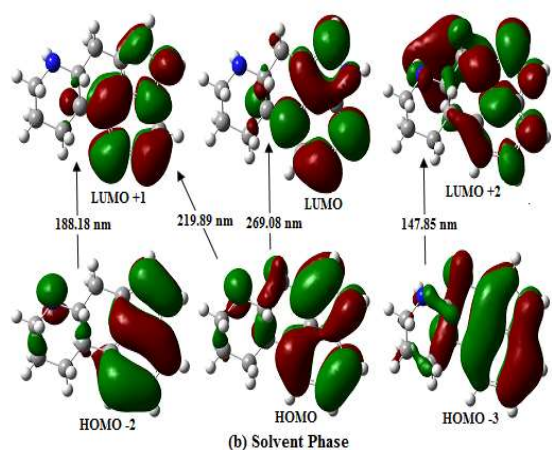


Figure 4: HOMO-LUMO plots (a) Gas phase and (b) Solvent phase.

Vibrational analysis

This molecule has 32 atoms and hence gives 90 ($3N-6$; N number of atoms) modes of vibration. All the modes are both Raman and IR active. Some selected fundamental vibrational wavenumbers, intensities and the PED distribution obtained by HF and DFT methods with 6-311++G (d,p) basis set are given in Table 2. Details of vibrational wavenumbers and their PED is listed in Table S (supplementary Material). The wavenumbers obtained by DFT have smaller values than that by HF methods due to inclusion of electron correlations in it. The assignments are given as per the internal coordinate system recommended by Pulay *et al.* 1979³¹ using DFT. The theoretical wavenumbers, generally higher than the observed values, were corrected by a scaling factor given by Yoshida *et al.*³⁷ using the expression: $\nu_{\text{obs}} = (1.0087 - 0.0000163 \nu_{\text{cal}}) \nu_{\text{cal}}$. The Raman scattering cross-sections, $\partial\sigma_j/\partial\nu$, which are proportional to the Raman intensities

may be calculated from the Raman scattering amplitude and predicted wavenumbers for each normal modes using the relationship^{38,39}:

$$\frac{\partial\sigma_j}{\partial\Omega} = \left(\frac{2^4\pi^4}{45}\right) \left(\frac{(v_0-v_j)^4}{1-\exp\left(-\frac{hcv_j}{kT}\right)}\right) \left(\frac{h}{8\pi^2cv_j}\right) S_j$$

where S_j and v_j are the scattering activities and the predicted wavenumbers (in cm^{-1}), respectively of the j^{th} normal mode, v_0 is the Raman exciting wavenumber (in cm^{-1}), and h , c and k are universal constants. The FT-IR and Raman spectra are given in figures 5 and 6, respectively. The assignments of some prominent modes are discussed below:

Ring R1 vibrations

The N-H stretching absorption generally occurred in the region $3500-3220 \text{ cm}^{-1}$.⁴⁰ In this study, pure N-H stretching vibration was calculated at 3320 cm^{-1} with IR intensity 0.4759 and Raman activity 73.268 units, respectively. The mixed in-plane-bending was calculated at 1473 cm^{-1} .

The CH_2 asymmetric stretching vibration was calculated at 2937, 2930 and 2919 cm^{-1} . The symmetric stretching of the mode was predicted at 2889, 2886 and 2878 cm^{-1} . The wagging vibration was predicted at 1371, 1366 and 1351 cm^{-1} in the scaled DFT. The twisting vibrations were calculated at 1308, 1274 and 1194 cm^{-1} . The C4-H /C3-H stretching vibrations calculated at $2877/2833 \text{ cm}^{-1}$ have intensities 42.173/11.63 units in the IR spectra and 115.142/106.889 units in the Raman spectra. The highly mixed rocking vibrations related to these modes were calculated at 1345 and 1338 cm^{-1} .

Table 2. Selected calculated IR and Raman wave numbers (cm^{-1}) of ergoline.

| DFT | | HF | Raman activity ($\text{\AA}^4/\text{a.m. u.}$) | IR intens. ($\text{km}\cdot\text{mol}^{-1}$) | Potient energy distribution (PED%) |
|-----------|--------|--------|--|--|------------------------------------|
| Unscale d | Scaled | scaled | | | |
| 3676 | 3488 | 3703 | 176.582 | 80.067 | R3[v(NH)](100) |
| 3488 | 3320 | 3540 | 73.268 | 0.4759 | R1[v(NH)]100) |
| 3241 | 3098 | 3233 | 134.288 | 1.289 | R3[v(CH)](98) |
| 3185 | 3047 | 3194 | 269.652 | 23.062 | R4[v(CH)](100) |
| 3174 | 3037 | 3183 | 83.815 | 19.732 | R4[v(CH)](97) |

| | | | | | |
|------|------|------|---------|--------|---|
| 3161 | 3026 | 3165 | 63.869 | 1.458 | R4[$\nu(\text{CH})$](86) |
| 3069 | 2942 | 3082 | 112.236 | 23.105 | R2[$\nu_a(\text{CH}_2)$](96) |
| 3063 | 2937 | 3069 | 156.889 | 56.344 | R1[$\nu_a(\text{CH}_2)$](97) |
| 3056 | 2930 | 3064 | 50.957 | 37.593 | R1[$\nu_a(\text{CH}_2)$](96) |
| 3044 | 2919 | 3050 | 231.545 | 81.004 | R1[$\nu_a(\text{CH}_2)$](96) |
| 3010 | 2889 | 3025 | 71.513 | 51.843 | R1[$\nu_s(\text{CH}_2)$](86)+ $\nu(\text{C}_4\text{H})$ (8)] |
| 3007 | 2886 | 3020 | 33.024 | 27.533 | R1[$\nu_s(\text{CH}_2)$](93) |
| 2999 | 2878 | 3015 | 127.905 | 14.386 | R1[$\nu_s(\text{CH}_2)$](60)+ $\nu(\text{C}_4\text{H})$ (27)] |
| 2997 | 2877 | 3005 | 115.142 | 42.173 | R1[$\nu(\text{C}_4\text{H})$](45)+ $\nu_s(\text{CH}_2)$](43)] |
| 2978 | 2859 | 3002 | 144.925 | 18.564 | R2[$\nu_s(\text{CH}_2)$](71)+R1[$\nu(\text{C}_4\text{H})$](15) |
| 2949 | 2833 | 2976 | 106.889 | 11.763 | R1[$\nu(\text{C}_3\text{H})$](90) |
| 1651 | 1621 | 1764 | 27.684 | 6.155 | R4[(CC)(45)+ δ_{trig} (26)]+R2[δ_{trig}](15)+R3[$\nu(\text{C}=\text{C})$](12) |
| 1633 | 1604 | 2976 | 27.304 | 12.490 | R3[δ'_{ring}](33)+ $\nu(\text{CC})$ (12)]+R2[$\nu(\text{CC})$](33)+R4[δ_a](14 |
| 1590 | 1563 | 1701 | 72.071 | 8.583 | R3[$\nu(\text{C}=\text{C})$](48)+ δ'_{ring} (18)]+R4[δ'_a](14)+R2[δ_{trig}](6) |
| 1535 | 1510 | 1632 | 3.843 | 3.897 | R4[$\delta_{\text{in}}(\text{CH})$](48)+R3[$\delta_{\text{in}}(\text{NH})$](23)+R2[$\nu(\text{CC})$](19) |
| 1499 | 1475 | 1608 | 1.483 | 3.198 | R1[$\delta(\text{CH}_2)$](68)+ $\rho(\text{CH}_2)$](19)] |
| 1496 | 1473 | 1597 | 0.777 | 13.512 | R1[$\delta_{\text{in}}(\text{NH})$](46)+ $\rho(\text{C}_4\text{H})$ (23)+ $\delta(\text{CH}_2)$](19)] |
| 1487 | 1464 | 1585 | 14.838 | 3.032 | R1[$\delta(\text{CH}_2)$](27)+ δ_{trig} (14)]+R2[$\delta(\text{CH}_2)$](22)+ δ'_{ring} (15)] |
| 1485 | 1462 | 1581 | 6.954 | 5.071 | R2[$\delta(\text{CH}_2)$](46)+R1[$\delta(\text{CH}_2)$](34) |
| 1479 | 1456 | 1577 | 12.824 | 17.448 | R1[$\delta(\text{CH}_2)$](37)+ $\delta_{\text{in}}(\text{NH})$ (30)+ $\rho(\text{CH}_2)$](11)] |
| 1458 | 1436 | 1551 | 11.749 | 21.852 | R4[δ'_a](30)+ $\delta_{\text{in}}(\text{CH})$ (21)]+R3[$\nu(\text{CC})$](26)+R2[δ_{trig}](11) |
| 1444 | 1423 | 1538 | 60.106 | 10.565 | R3[$\delta_{\text{in}}(\text{NH})$](25)+ $\nu(\text{NC})$ (17)+R2[$\nu(\text{CC})$](30)+R4[δ'_a](15) |
| 1395 | 1375 | 1499 | 39.051 | 2.082 | R3[δ'_{ring}](21)+ $\nu(\text{CC})$ (14)]+R2[$\nu(\text{CC})$](27)+R4[δ_{trig}](15) |
| 1390 | 1371 | 1494 | 16.434 | 1.897 | R1[ω](18) γ (8)(CH_2)+ $\rho(\text{CH})$ (15)]+R2[$\omega(\text{CH}_2)$](18)+ $\nu(\text{CC})$ (11)] |
| 1385 | 1366 | 1479 | 0.763 | 1.014 | R1[$\omega(\text{CH}_2)$](42)+ δ_{trig} (23)]+R2[$\rho(\text{C}_3\text{H})$](12)+ δ'_a (9)] |
| 1373 | 1354 | 1475 | 4.532 | 6.295 | R1[$\omega(\text{CH}_2)$](35)+ δ_{trig} (14)]+R4[$\nu(\text{CC})$](21) |
| 1372 | 1353 | 1472 | 1.715 | 18.654 | R2[$\omega(\text{CH}_2)$](24)+ $\nu(\text{CC})$ (8)]+R1[$\omega(\text{CH}_2)$](31)+R3[$\nu(\text{C}=\text{C})$](18) |
| 1356 | 1338 | 1426 | 7.158 | 9.850 | R1[$\rho'(\text{CH})$](39)+ $\gamma(\text{CH}_2)$ (15)]+R2[$\gamma(\text{CH}_2)$](17)+ δ'_a (15)] |
| 1309 | 1292 | 1377 | 18.185 | 15.722 | R2[$\nu(\text{CC})$](35)+ δ'_a (5)]+R4[$\delta_{\text{in}}(\text{CH})$](14)+ δ_{trig} (12)]+R3[$\nu(\text{NC})$](23) |
| 1290 | 1274 | 1369 | 5.039 | 1.007 | R1[$\gamma(\text{CH}_2)$](46)+R2[δ'_a](20)+R2[$\gamma(\text{CH}_2)$](20) |
| 1285 | 1269 | 1367 | 4.722 | 1.603 | R4[δ_{trig}](44)+ $\nu(\text{CC})$ (5)]+R3[δ'_{ring}](20)+ $\nu(\text{CC})$ (9)] |
| 1253 | 1238 | 1332 | 2.294 | 17.712 | R2[δ'_a](40)+ δ_{trig} (7)]+R4[δ_{trig}](28)+R3[δ'_{ring}](10) |
| 1240 | 1226 | 1324 | 9.112 | 7.109 | R1[δ_{trig}](44)+R2[δ_{trig}](22)+R3[$\nu(\text{NC})$](9)+ $\nu(\text{CC})$ (5)] |
| 1234 | 1220 | 1313 | 14.237 | 3.165 | R4[δ_{trig}](44)+R3[$\nu(\text{CC})$](13)+ δ'_{ring} (12)+ $\nu(\text{NC})$ (12)]+R2[δ'_a](8) |

| | | | | | |
|------|------|------|-------|--------|---|
| 1207 | 1194 | 1288 | 0.981 | 0.981 | R2[γ (CH ₂)](53)+R1[γ (CH ₂)](11)+R3[ν (NC)](8) |
| 1183 | 1170 | 1261 | 4.232 | 2.784 | R4[δ_{in} (CH)(21)+ ν (CC)(14)+ δ_{trig} (13)]+R2[ν (CC)](13) |
| 1178 | 1166 | 1238 | 2.871 | 0.0854 | R2[δ_{trig} (50)+ δ'_a (12)]+R1[δ_{trig} (15)]+R3[δ'_{ring}](6) |
| 1147 | 1136 | 1223 | 5.845 | 33.362 | R2[puck](48)+ R2[ν (CC)](7)+R1[δ'_a](25) |
| 1132 | 1121 | 1211 | 6.585 | 14.422 | R1[δ_{trig}](44)+R2[δ_{trig} (17)+ δ'_a (15)]+R4[δ'_{ring}](8) |
| 1116 | 1105 | 1100 | 1.220 | 13.348 | R4[δ_{trig}](45)+R3[δ'_{ring}](27)+R2[δ_{trig}](13) |
| 1101 | 1091 | 1174 | 3.647 | 7.290 | R2[δ_{trig} (42)+ δ'_a (17)]+R3[ν (NC)](15)+R1[ν (CC)](8) |
| 1088 | 1078 | 1148 | 1.298 | 22.097 | R3[δ'_{ring} (37)+ ν (NC)(33)]+R2[δ'_a](10) |
| 1069 | 1060 | 1139 | 2.787 | 3.401 | R4[δ_{trig}](39)+R2[ν (CC)](21)+R3[δ'_{ring}](11) |
| 920 | 914 | 986 | 6.390 | 4.841 | R4[δ_{trig} (57)+ δ_a (8)]+R2[δ_{trig}](12)+R3[δ'_{ring}](12) |
| 897 | 892 | 978 | 1.551 | 8.033 | R2[puck(48)+ δ'_a (11)]+R3[δ'_{ring}](18)+R4[puck](7) |
| 881 | 876 | 958 | 0.728 | 3.817 | R2[δ_{trig} (37)+ δ'_a (15)]+R4[oop(CH)](27) |
| 862 | 857 | 919 | 6.980 | 1.149 | R4[puck(56)+ δ'_a (24)] |
| 849 | 845 | 916 | 1.246 | 1.093 | R4[δ_{trig} (41)+ δ'_a (37)]+R3[δ'_{ring}](15) |
| 821 | 817 | 903 | 2.313 | 32.478 | R2[δ'_a (42)+ ν (CC)(6)]+R1[ν (CC)](38) |
| 730 | 728 | 788 | 3.384 | 42.321 | R4[δ_a (20)+ δ_{trig} (15)]+ R3[δ_{ring}](26)+R2[δ'_a](25) |
| 635 | 634 | 684 | 2.487 | 7.073 | R2[δ_{trig}](39)+R3[δ'_{ring}](35)+R4[δ_a](19) |
| 615 | 614 | 665 | 3.672 | 1.377 | R3[δ'_{ring}](69)+R4[δ_a](12)+R2[δ_{trig}](9) |
| 602 | 601 | 649 | 8.028 | 0.419 | R4[δ'_a](58)+R2[δ_{trig} (19)+ δ'_a (8)]+R3[δ'_{ring}](5) |
| 595 | 594 | 637 | 2.146 | 1.526 | R3[δ'_{ring}](46)+R4[δ'_a](43) |
| 551 | 551 | 589 | 6.959 | 1.725 | R4[δ'_a](51)+R2[δ'_a (36)+ δ_{trig} (5)] |

Proposed assignments and potential energy distribution (PED) for vibrational normal modes.

Types of vibration: ν , stretching; δ , deformation (bending),

scissoring; oop, out-of-plane bending; ω , wagging; γ , twisting; ρ , rocking; τ , torsion.

Potential energy distribution (contribution ≥ 5).

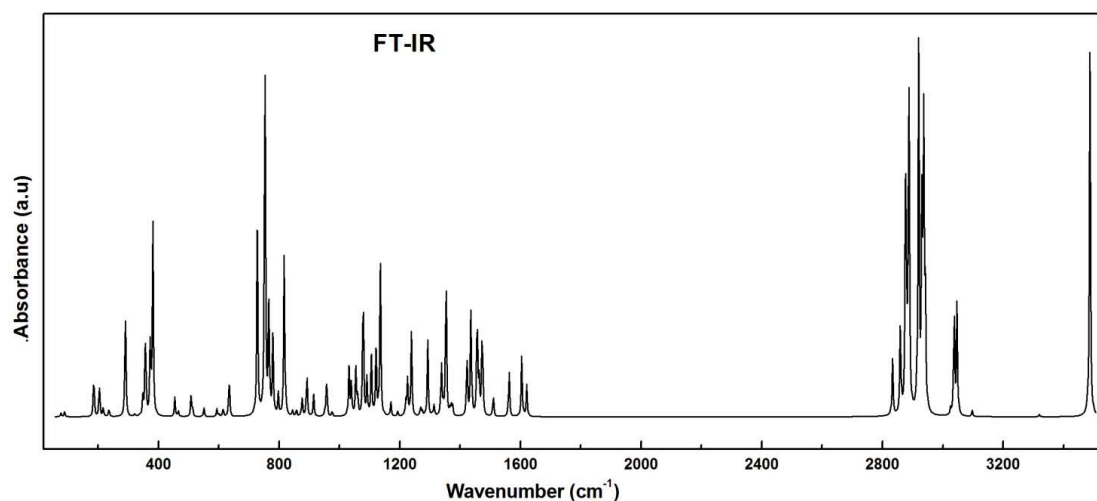


Figure 5: Calculated FT-IR spectrum between the ranges 200-3500 cm⁻¹.

Ring R2 vibrations

The asymmetric stretching of CH₂ predicted at 2942 cm⁻¹ has intensity 23.105/112.236 units in the IR/Raman spectra. The symmetric stretching was calculated at 2859 cm⁻¹. The deformation, wagging and twisting mode vibrations were calculated at 1462, 1353 and 1194 cm⁻¹, respectively.

Highly mixed C-C ring stretching was predicted at 1292 cm⁻¹. The ring deformation and puckering vibrations are below 1200 cm⁻¹ in the mixed modes.

Ring R3 vibrations

The N-H stretching vibration was calculated at 348 cm⁻¹. This pure mode (100% contribution in PED) has intensity 176.582 units in the Raman spectra and 80.067 units in the IR spectrum. The in-plane vibration of this mode was predicted at 1423 cm⁻¹, while out-of-plane vibration was at 381 cm⁻¹ as listed in the Table 2.

The frequency of ring deformation vibration was calculated at 1604 cm⁻¹. The C=C stretching predicted at 1563 cm⁻¹ is a mixed mode. The mixed NC stretching is below 1430 cm⁻¹.

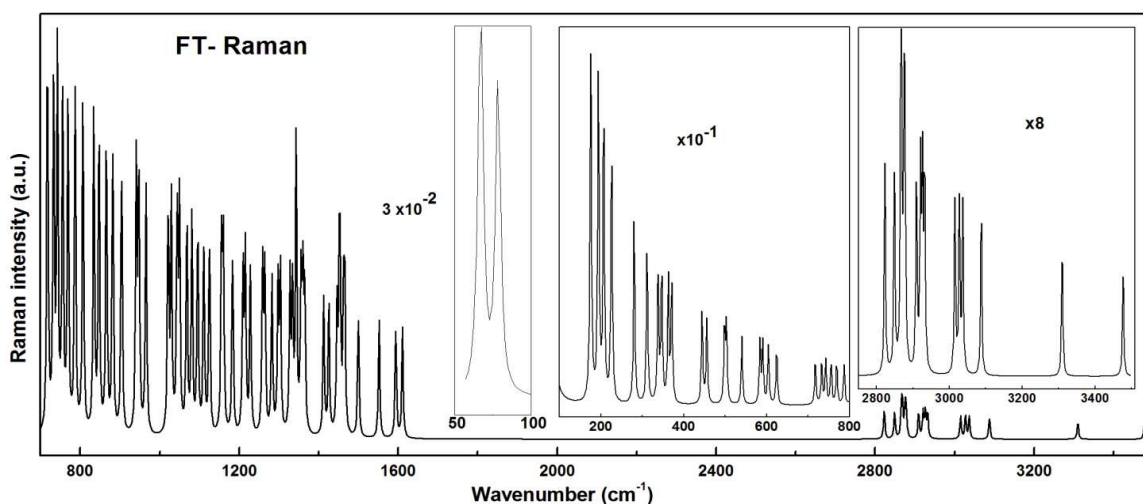


Figure 6: Calculated FT-Raman spectrum between the ranges 50-3500 cm⁻¹ (intensity magnitudes shown in the insets are modified as: 50-100 cm⁻¹ by 3x10⁻², 100-800 cm⁻¹ by 10⁻¹ and 2750-3500 cm⁻¹ by 8 times, respectively).

Ring R4 vibrations

One can easily expect three modes of stretching vibration related to each CH moiety. In this study, the C-H stretching vibration was predicted at 3047, 3037 and 3026 cm⁻¹. The in-plane deformation was calculated at 1510 cm⁻¹. The out-of-plane vibration of this mode is 951 cm⁻¹. In the ring, C-C stretching and deformation vibrations were predicted at 1621 and 1436 cm⁻¹, respectively. The trigonal ring deformations were calculated below 1269 cm⁻¹.

distribution has been analyzed by *ab initio* and DFT methods. Information about the shape, size, charge density and the structure-activity relation have been analyzed by mapping MEP. The large value of the band gap obtained from the HOMO and LUMO analysis supports that the title molecule is stable; however, there occur intramolecular charge transfer (ICT) within the molecule. This study could be a better starting for further investigations on structure-activity relation of the compound.

CONCLUSION

The investigated molecule has 90 modes of vibrations. All the modes are both IR and Raman active. In this study, the vibrational wavenumbers along with their potential energy

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